

DEM Modeling: Lecture 13

3D Rotations

3D Rotations

- Newton's 2nd Law:

$$\frac{d}{dt}(\underline{\underline{I}} \cdot \underline{\omega})^G = \mathbf{T}^G$$

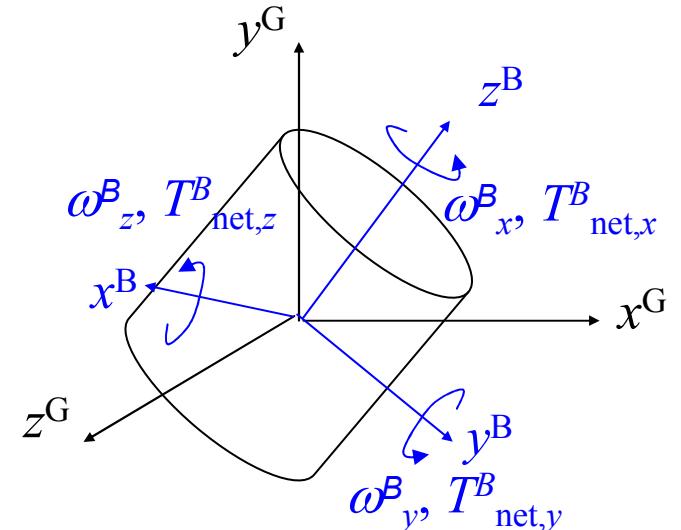
- Easiest to put the rotational eqns of motion in a body-fixed frame of reference (FOR) aligned with the particle's principle axes
 - in a body-fixed FOR, the moments of inertia don't change due to changes in orientation
 - aka Euler's rotational eqns of motion

$$\dot{\underline{\omega}}^B = \underline{\underline{I}}^{-1} \cdot [\mathbf{T}^B - \underline{\omega}^B \times (\underline{\underline{I}} \cdot \underline{\omega}^B)]$$

$$\dot{\omega}_x^B = \frac{1}{I_{xx}} [T_x^B + \omega_y^B \omega_z^B (I_{yy} - I_{zz})]$$

$$\dot{\omega}_y^B = \frac{1}{I_{yy}} [T_y^B + \omega_z^B \omega_x^B (I_{zz} - I_{xx})]$$

$$\dot{\omega}_z^B = \frac{1}{I_{zz}} [T_z^B + \omega_x^B \omega_y^B (I_{xx} - I_{yy})]$$

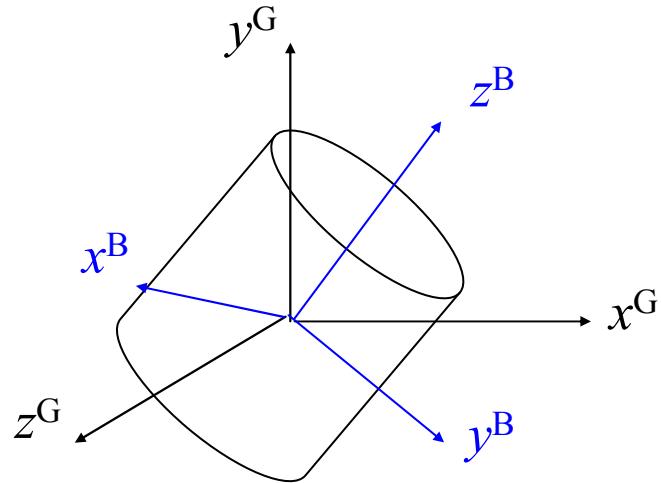


where

- I_{xx} , I_{yy} , and I_{zz} are the principle moments of inertia for the particle
- ω^B is the rotational speed of the particle in a body-fixed FOR
- \mathbf{T}^B is the net torque acting on the particle in the body-fixed FOR

3D Rotations...

- How should we describe the orientation of an object in 3D space?

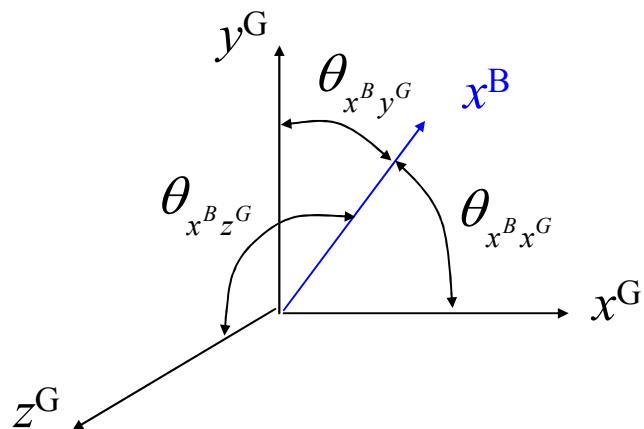


superscript “B” = body-fixed FOR
superscript “G” = global FOR

- Three common methods
 - direction cosines
 - Euler angles
 - quaternions

Direction Cosines

- Use the cosines of the angles that the body-fixed axes make with respect to the global axes to describe the orientation of the object



$$\hat{\mathbf{e}}_x^B = \hat{\mathbf{e}}_x^G \cos \theta_{x^B x^G} + \hat{\mathbf{e}}_y^G \cos \theta_{x^B y^G} + \hat{\mathbf{e}}_z^G \cos \theta_{x^B z^G}$$

$$\begin{Bmatrix} \hat{\mathbf{e}}_x^B \\ \hat{\mathbf{e}}_y^B \\ \hat{\mathbf{e}}_z^B \end{Bmatrix} = \underbrace{\begin{bmatrix} \cos \theta_{x^B x^G} & \cos \theta_{x^B y^G} & \cos \theta_{x^B z^G} \\ \cos \theta_{y^B x^G} & \cos \theta_{y^B y^G} & \cos \theta_{y^B z^G} \\ \cos \theta_{z^B x^G} & \cos \theta_{z^B y^G} & \cos \theta_{z^B z^G} \end{bmatrix}}_{=R^{BG}} \begin{Bmatrix} \hat{\mathbf{e}}_x^G \\ \hat{\mathbf{e}}_y^G \\ \hat{\mathbf{e}}_z^G \end{Bmatrix}$$

$R^{BG} \equiv$ rotation matrix from the Global to the Body FOR

Direction Cosines...

- The rotation matrix R^{BG} is orthonormal
 - the basis vectors forming the rows and columns of the matrix are mutually perpendicular
 - the magnitude of the basis vectors formed by these rows and columns have a magnitude of one
 - $\Rightarrow R^{GB} = (R^{BG})^{-1} = (R^{BG})^T$

$$\begin{Bmatrix} \hat{\mathbf{e}}_x^G \\ \hat{\mathbf{e}}_y^G \\ \hat{\mathbf{e}}_z^G \end{Bmatrix} = R^{GB} \begin{Bmatrix} \hat{\mathbf{e}}_x^B \\ \hat{\mathbf{e}}_y^B \\ \hat{\mathbf{e}}_z^B \end{Bmatrix}$$

- although the matrix R^{GB} has nine terms, only three are independent

Direction Cosines...

- *Example*

- The unit vectors for the body-fixed FOR written in terms of the global FOR are:

$$\hat{\mathbf{e}}_x^B = \hat{\mathbf{e}}_y^G$$

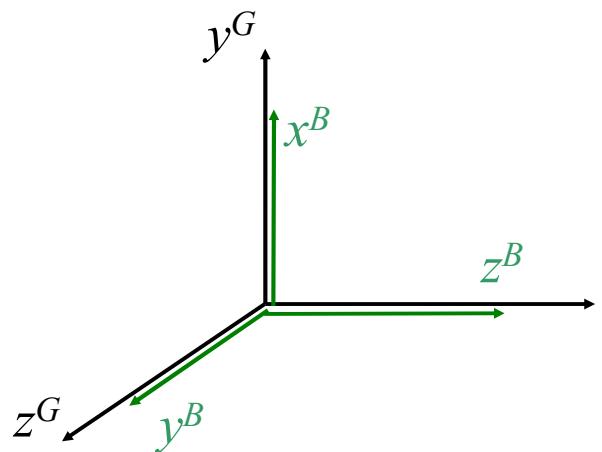
$$\hat{\mathbf{e}}_y^B = \hat{\mathbf{e}}_z^G$$

$$\hat{\mathbf{e}}_z^B = \hat{\mathbf{e}}_x^G$$

- Sketch both the body-fixed and global FOR axes. Let the global and body-fixed FORs share the same origin.
 - Write the rotation matrix for the transformation from the body-fixed FOR to the global FOR.
 - Express the vector $\mathbf{v}^B = (0, 1, 0)$ which is given in the body-fixed FOR in terms of the global FOR using the rotation matrix determined in the previous part.

Direction Cosines...

- Solution*



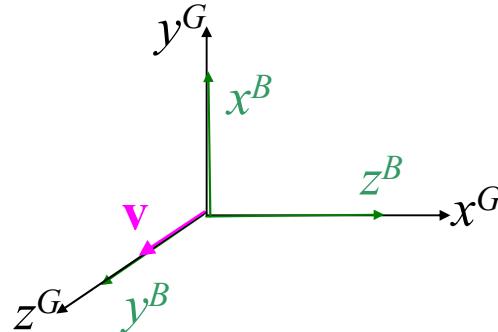
$$R^{BG} = \begin{bmatrix} \cos \theta_{x^B x^G} & \cos \theta_{x^B y^G} & \cos \theta_{x^B z^G} \\ \cos \theta_{y^B x^G} & \cos \theta_{y^B y^G} & \cos \theta_{y^B z^G} \\ \cos \theta_{z^B x^G} & \cos \theta_{z^B y^G} & \cos \theta_{z^B z^G} \end{bmatrix} = \begin{bmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & 0 & 0 \end{bmatrix}$$

$$R^{GB} = (R^{BG})^{-1} = (R^{BG})^T = \begin{bmatrix} 0 & 0 & 1 \\ 1 & 0 & 0 \\ 0 & 1 & 0 \end{bmatrix}$$

$$\mathbf{v}^G = R^{GB} \mathbf{v}^B$$

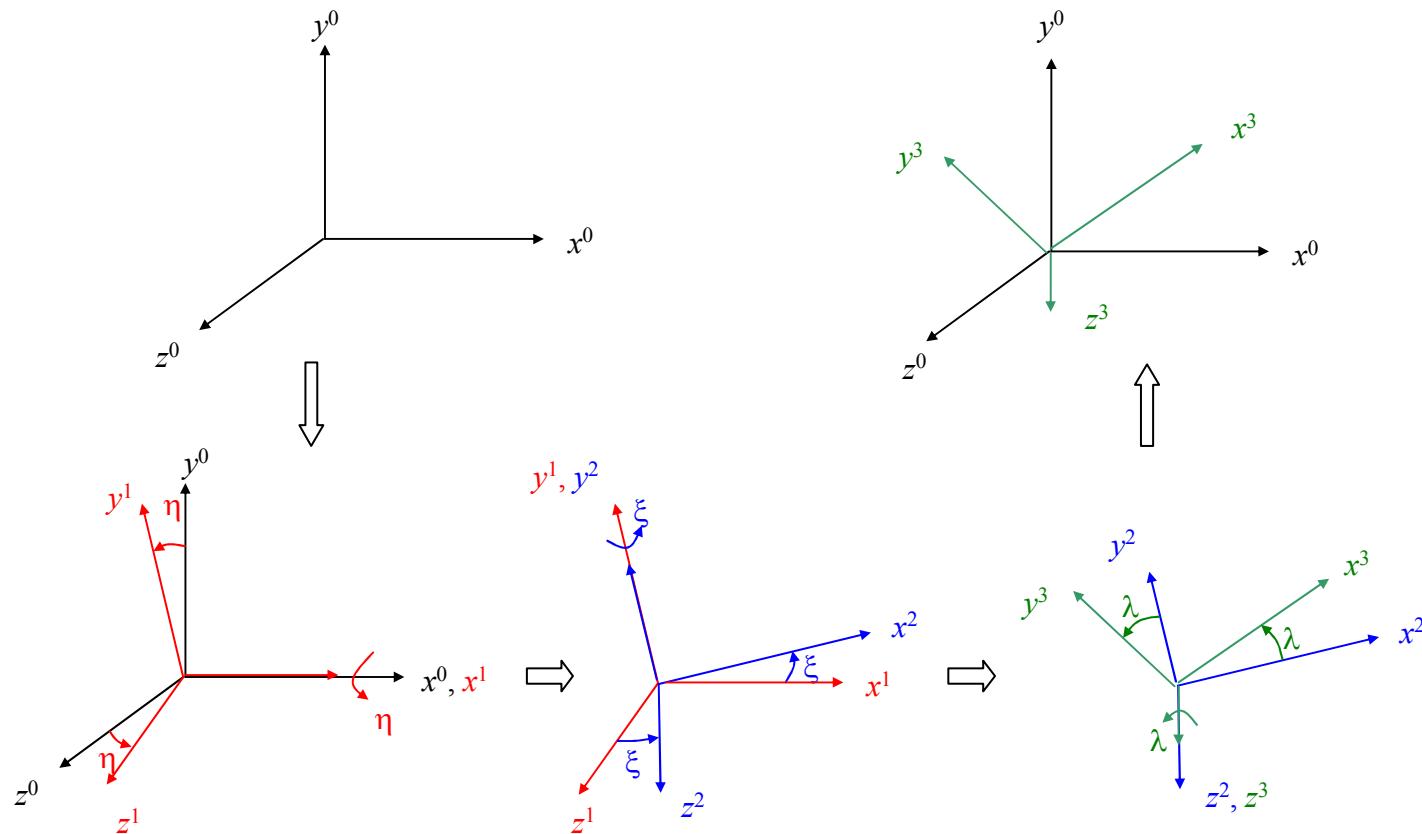
$$= \begin{bmatrix} 0 & 0 & 1 \\ 1 & 0 & 0 \\ 0 & 1 & 0 \end{bmatrix} \begin{Bmatrix} 0 \\ 1 \\ 0 \end{Bmatrix}$$

$$\therefore \mathbf{v}^G = \begin{Bmatrix} 0 \\ 0 \\ 1 \end{Bmatrix}$$



Euler Angles

- Decompose the orientation into three successive rotations about various coordinate axes (e.g. the 1-2-3 sequence shown below)



Euler Angles...

- The full sequence of rotations may be expressed as a single rotation matrix

$$\mathbf{a}^1 = R^{10} \mathbf{a}^0 = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos \eta & \sin \eta \\ 0 & -\sin \eta & \cos \eta \end{bmatrix} \mathbf{a}^0$$

$$\mathbf{a}^3 = R^{30} \mathbf{a}^0 = R^{32} R^{21} R^{10} \mathbf{a}^0$$

$$\mathbf{a}^2 = R^{21} \mathbf{a}^1 = \begin{bmatrix} \cos \xi & 0 & -\sin \xi \\ 0 & 1 & 0 \\ \sin \xi & 0 & \cos \xi \end{bmatrix} \mathbf{a}^1$$

$$R^{30} = \begin{bmatrix} c\xi c\lambda & c\eta s\lambda + s\eta s\xi c\lambda & s\eta s\lambda - c\eta s\xi c\lambda \\ -c\xi s\lambda & c\eta c\lambda - s\eta s\xi s\lambda & s\eta c\lambda + c\eta s\xi s\lambda \\ s\xi & -s\eta c\xi & c\eta c\xi \end{bmatrix}$$

$$\mathbf{a}^3 = R^{32} \mathbf{a}^2 = \begin{bmatrix} \cos \lambda & \sin \lambda & 0 \\ -\sin \lambda & \cos \lambda & 0 \\ 0 & 0 & 1 \end{bmatrix} \mathbf{a}^2$$

Euler Angles...

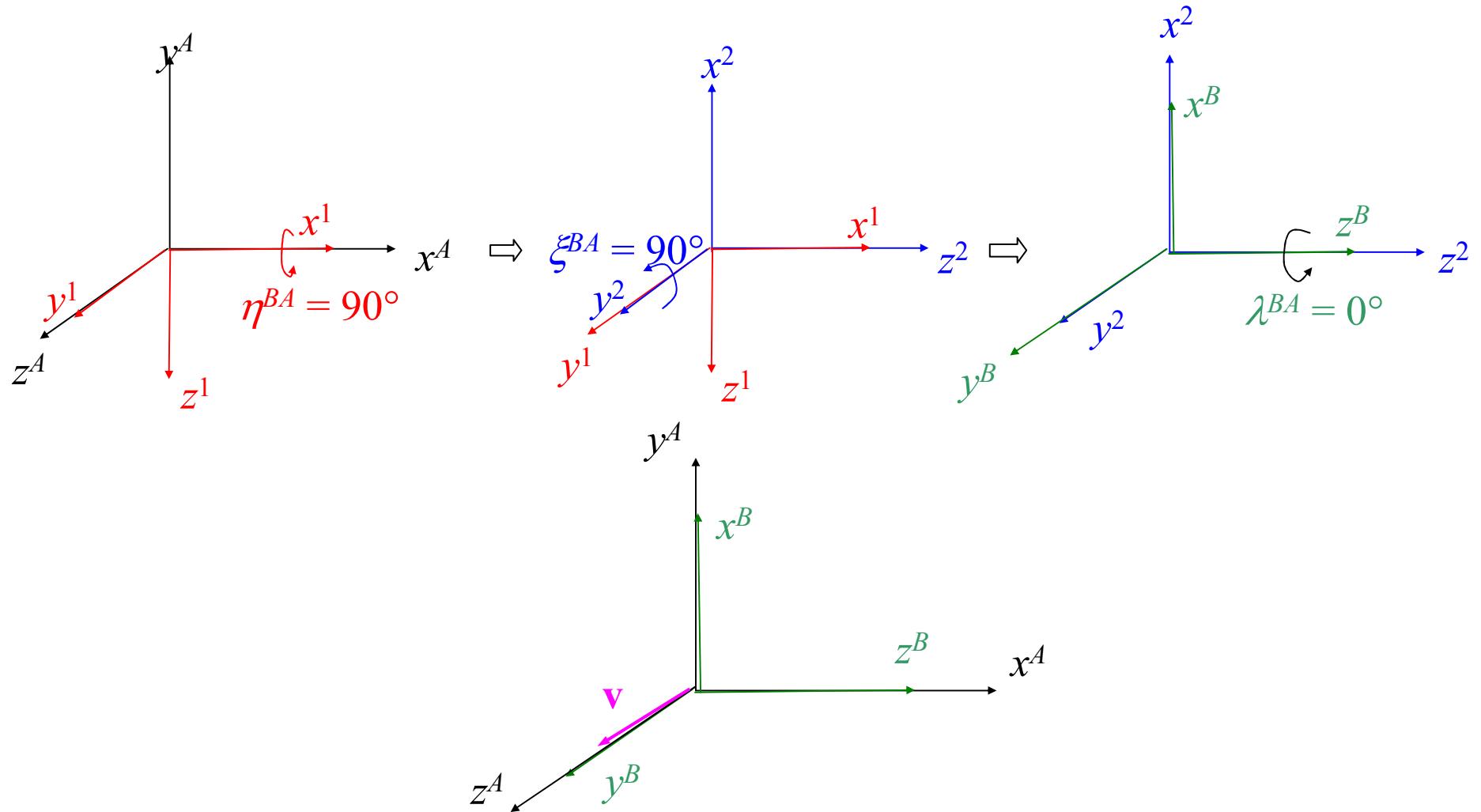
- There are twelve possible Euler angle sequences that could be used:.. 1-2-1 (e.g. $x^0y^1x^2$), 1-3-1, 2-1-2, 2-3-2, 3-1-3, 3-2-3, 1-2-3, 2-3-1, 3-1-2, 3-2-1, 2-1-3, and 1-3-2.
- The rotation matrix is orthonormal.
- The Euler angles used to achieve a final orientation are not unique. For example, the rotation matrix is identical for $(\eta, \xi, \lambda) = (135^\circ, 135^\circ, 135^\circ)$ and $(\eta, \xi, \lambda) = (-45^\circ, 45^\circ, -45^\circ)$.

Euler Angles...

- *Example*
 - A vector $\mathbf{v}^B = (0, 1, 0)$ is expressed in FOR B. The Euler angles for converting from FOR A to FOR B are $(\eta, \xi, \lambda)^{BA} = (90^\circ, 90^\circ, 0^\circ)$. Sketch FORs A and B and the vector \mathbf{v} . Find vector \mathbf{v} 's components in FOR A.

Euler Angles...

- *Solution*



Euler Angles...

- *Solution...*

$$\mathbf{v}^B = R^{BA} \mathbf{v}^A \Rightarrow \mathbf{v}^A = (R^{BA})^{-1} \mathbf{v}^B = (R^{BA})^T \mathbf{v}^B$$

$$R^{BA} = \begin{bmatrix} c\xi^{BA}c\lambda^{BA} & c\eta^{BA}s\lambda^{BA} + s\eta^{BA}s\xi^{BA}c\lambda^{BA} & s\eta^{BA}s\lambda^{BA} - c\eta^{BA}s\xi^{BA}c\lambda^{BA} \\ -c\xi^{BA}s\lambda^{BA} & c\eta^{BA}c\lambda^{BA} - s\eta^{BA}s\xi^{BA}s\lambda^{BA} & s\eta^{BA}c\lambda^{BA} + c\eta^{BA}s\xi^{BA}s\lambda^{BA} \\ s\xi^{BA} & -s\eta^{BA}c\xi^{BA} & c\eta^{BA}c\xi^{BA} \end{bmatrix}$$

$$= \begin{bmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & 0 & 0 \end{bmatrix}$$

$$\begin{Bmatrix} v_x \\ v_y \\ v_z \end{Bmatrix}^A = \underbrace{\begin{bmatrix} 0 & 0 & 1 \\ 1 & 0 & 0 \\ 0 & 1 & 0 \end{bmatrix}}_{= (R^{BA})^T} \begin{Bmatrix} 0 \\ 1 \\ 0 \end{Bmatrix}^B$$

$$\therefore \mathbf{v}^A = \begin{Bmatrix} v_x \\ v_y \\ v_z \end{Bmatrix}^A = \begin{Bmatrix} 0 \\ 0 \\ 1 \end{Bmatrix}$$

Euler Angles...

- Euler angle rate of change
 - consider a small rotation about each of the Euler angle sequence axes

$$\Delta\theta = \Delta\eta \hat{\mathbf{e}}_x^0 + \Delta\xi \hat{\mathbf{e}}_y^1 + \Delta\lambda \hat{\mathbf{e}}_z^2$$

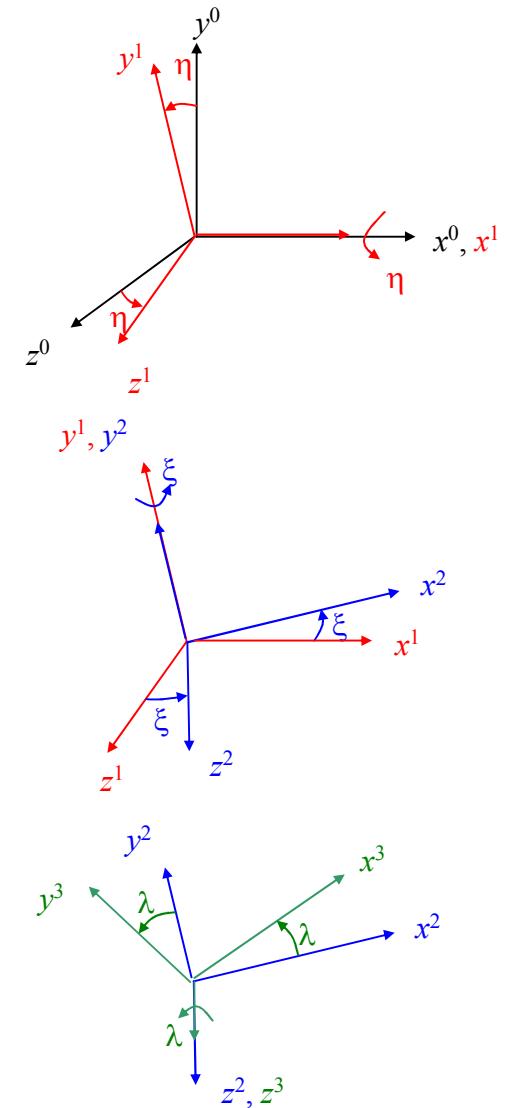
- express the rate at which the rotation, $\Delta\theta$, occurs in terms of the body-fixed FOR (frame 3)

$$\hat{\mathbf{e}}_z^2 = \hat{\mathbf{e}}_z^3$$

$$\hat{\mathbf{e}}_y^1 = \hat{\mathbf{e}}_y^2 = \sin \lambda \hat{\mathbf{e}}_x^3 + \cos \lambda \hat{\mathbf{e}}_y^3$$

$$\hat{\mathbf{e}}_x^0 = \hat{\mathbf{e}}_x^1 = \cos \xi \hat{\mathbf{e}}_x^2 + \sin \xi \hat{\mathbf{e}}_z^2$$

$$\hat{\mathbf{e}}_x^2 = \cos \lambda \hat{\mathbf{e}}_x^3 - \sin \lambda \hat{\mathbf{e}}_y^3$$



Euler Angles...

- Euler angle rate of change...

$$\Delta\theta^3 = (\Delta\eta \cos\xi \cos\lambda + \Delta\xi \sin\lambda) \hat{\mathbf{e}}_x^3 + (-\Delta\eta \cos\xi \sin\lambda + \Delta\xi \cos\lambda) \hat{\mathbf{e}}_y^3 + (\Delta\eta \sin\xi + \Delta\lambda) \hat{\mathbf{e}}_z^3$$

$$\underbrace{\frac{d\theta^3}{dt}}_{=\omega^3} = \left(\underbrace{\frac{d\eta}{dt}}_{=\dot{\eta}} \cos\xi \cos\lambda + \underbrace{\frac{d\xi}{dt}}_{=\dot{\xi}} \sin\lambda \right) \hat{\mathbf{e}}_x^3 + \left(-\underbrace{\frac{d\eta}{dt}}_{=\dot{\eta}} \cos\xi \sin\lambda + \underbrace{\frac{d\xi}{dt}}_{=\dot{\xi}} \cos\lambda \right) \hat{\mathbf{e}}_y^3 + \left(\underbrace{\frac{d\eta}{dt}}_{=\dot{\eta}} \sin\xi + \underbrace{\frac{d\lambda}{dt}}_{=\dot{\lambda}} \right) \hat{\mathbf{e}}_z^3$$

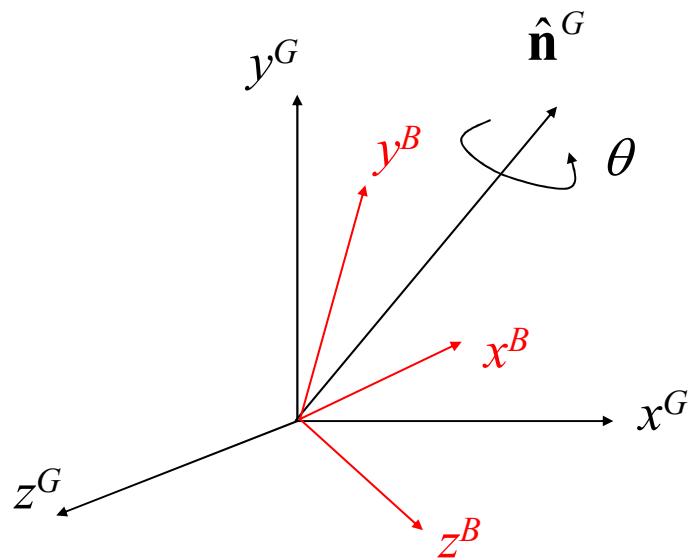
$$\begin{Bmatrix} \omega_x^3 \\ \omega_y^3 \\ \omega_z^3 \end{Bmatrix} = \begin{bmatrix} \cos\xi \cos\lambda & \sin\lambda & 0 \\ -\cos\xi \sin\lambda & \cos\lambda & 0 \\ \sin\xi & 0 & 1 \end{bmatrix} \begin{Bmatrix} \dot{\eta} \\ \dot{\xi} \\ \dot{\lambda} \end{Bmatrix}$$

- if $\xi \rightarrow \pi/2$ or $3\pi/2$, then $\cos\xi \rightarrow 0$
- The Euler angle formulation is not well posed for all Euler angles.
- Other Euler angle sequences (e.g. 1-2-1) will also have singularities at particular angles.
- This behavior is sometimes referred to as “gimbal lock.”

$$\therefore \begin{Bmatrix} \dot{\eta} \\ \dot{\xi} \\ \dot{\lambda} \end{Bmatrix} = \begin{bmatrix} \cos\lambda / \cos\xi & -\sin\lambda / \cos\xi & 0 \\ \sin\lambda & \cos\lambda & 0 \\ -\tan\xi \cos\lambda & \tan\xi \sin\lambda & 1 \end{bmatrix} \begin{Bmatrix} \omega_x^3 \\ \omega_y^3 \\ \omega_z^3 \end{Bmatrix}$$

Quaternions

- Describe an orientation using a single rotation about a unit vector



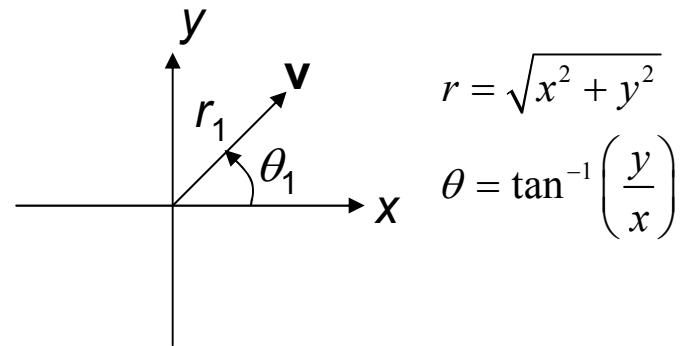
- Unlike direction cosines and Euler angles, a quaternion uses four quantities to describe the orientation: $(\theta, n_x^G, n_y^G, n_z^G)$. One of the quantities is not independent.
- An Euler angle rotation sequence can be thought of as a sequence of three quaternion rotations.

Quaternions...

- Consider a complex number

Let the 2D vector \mathbf{v} be represented by the complex number z_1

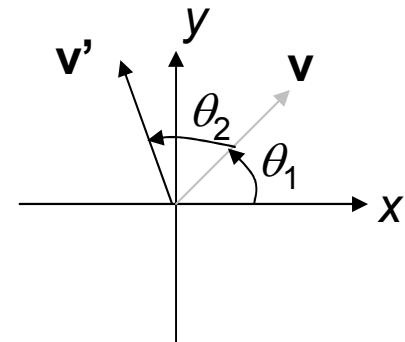
$$\mathbf{v} = z_1 = (x_1, y_1) = (r_1 \cos \theta_1, r_1 \sin \theta_1)$$



We can rotate the vector \mathbf{v} about the origin by an angle θ_2 by multiplying z_1 by the complex number z_2 as given below.

$$z_2 = (\cos \theta_2, \sin \theta_2)$$

$$\begin{aligned} \mathbf{v}' &= z_1 z_2 = (r_1 \cos \theta_1, r_1 \sin \theta_1)(\cos \theta_2, \sin \theta_2) \\ &= (r_1 \cos(\theta_1 + \theta_2), r_1 \sin(\theta_1 + \theta_2)) \end{aligned}$$



.: 2D rotations may be performed using complex numbers

Quaternions...

- (Unit) quaternions are “hypercomplex” numbers that can be used to perform 3D rotations
 - originally proposed by Hamilton (1843)

$$\hat{\mathbf{q}} = q_0 + q_1 i + q_2 j + q_3 k$$

$$= \left[\cos\left(\frac{\theta}{2}\right), \hat{\mathbf{n}} \sin\left(\frac{\theta}{2}\right) \right]$$

$$= \cos\left(\frac{\theta}{2}\right) + \sin\left(\frac{\theta}{2}\right)n_1 i + \sin\left(\frac{\theta}{2}\right)n_2 j + \sin\left(\frac{\theta}{2}\right)n_3 k$$

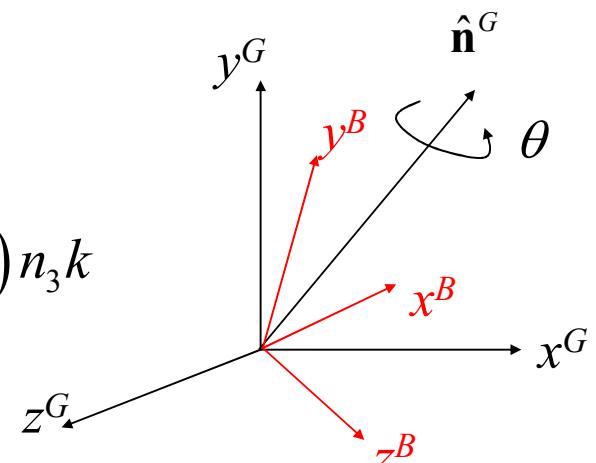
$$i^2 = j^2 = k^2 = -1$$

$$ij = -ji = k$$

$$jk = -kj = i$$

$$ki = -ik = j$$

$$i \\ k \circlearrowleft j \quad \quad i \\ k \circlearrowright j$$



Quaternions...

- Some handy quaternion properties

$$|\mathbf{q}| = \sqrt{q_0^2 + q_1^2 + q_2^2 + q_3^2}$$

$$\hat{\mathbf{q}} = \frac{\mathbf{q}}{|\mathbf{q}|}$$

$$\mathbf{a} + \mathbf{b} = \mathbf{b} + \mathbf{a}$$

$$\mathbf{a} + (\mathbf{b} + \mathbf{c}) = (\mathbf{a} + \mathbf{b}) + \mathbf{c}$$

$$\mathbf{a}(\mathbf{b}\mathbf{c}) = (\mathbf{a}\mathbf{b})\mathbf{c}$$

$$\mathbf{ab} \neq \mathbf{ba}, \text{ in general}$$

$$\begin{aligned}
 \mathbf{ab} &= [s, \mathbf{u}][t, \mathbf{v}] = [st - \mathbf{u} \cdot \mathbf{v}, s\mathbf{v} + t\mathbf{u} + \mathbf{u} \times \mathbf{v}] \\
 &= (a_0 + ia_1 + ja_2 + ka_3)(b_0 + ib_1 + jb_2 + kb_3) \\
 &= (a_0b_0 - a_1b_1 - a_2b_2 - a_3b_3) + i(a_0b_1 + a_1b_0 + a_2b_3 - a_3b_2) + \\
 &\quad j(a_0b_2 - a_1b_3 + a_2b_0 + a_3b_1) + k(a_0b_3 + a_1b_2 - a_2b_1 + a_3b_0)
 \end{aligned}$$

Quaternions...

- Some handy quaternion properties...

$$\mathbf{q}\mathbf{q}^{-1} = [1, \mathbf{0}]$$

if $\mathbf{q} = [s, \mathbf{v}] = (q_0, q_1, q_2, q_3)$, then $\mathbf{q}^{-1} = [s, -\mathbf{v}] = (q_0, -q_1, -q_2, -q_3)$

$$\mathbf{p}^{-1}\mathbf{q}^{-1} = (\mathbf{q}\mathbf{p})^{-1}$$

A vector \mathbf{v} rotated about unit vector $\hat{\mathbf{n}}$ by an angle θ is given by:

$$[0, \mathbf{v}'] = \hat{\mathbf{q}}[0, \mathbf{v}]\hat{\mathbf{q}}^{-1} = [0, \mathbf{v}\cos\theta + (\hat{\mathbf{n}} \times \mathbf{v})\sin\theta + (\mathbf{v} \cdot \hat{\mathbf{n}})\hat{\mathbf{n}}(1 - \cos\theta)]$$

Two successive rotations (first rotate using $\hat{\mathbf{q}}_1$ then rotate using $\hat{\mathbf{q}}_2$):

$$[0, \mathbf{v}''] = (\hat{\mathbf{q}}_2\hat{\mathbf{q}}_1)[0, \mathbf{v}](\hat{\mathbf{q}}_1^{-1}\hat{\mathbf{q}}_2^{-1}) \quad (\text{equivalent to } [0, \mathbf{v}''] = \hat{\mathbf{r}}[0, \mathbf{v}]\hat{\mathbf{r}}^{-1} \text{ where } \hat{\mathbf{r}} = \hat{\mathbf{q}}_2\hat{\mathbf{q}}_1)$$

Quaternions...

- Some handy quaternion properties...

$$[0, \mathbf{v}'] = \hat{\mathbf{q}} [0, \mathbf{v}] \hat{\mathbf{q}}^{-1} \Rightarrow \mathbf{v}' = R\mathbf{v}$$

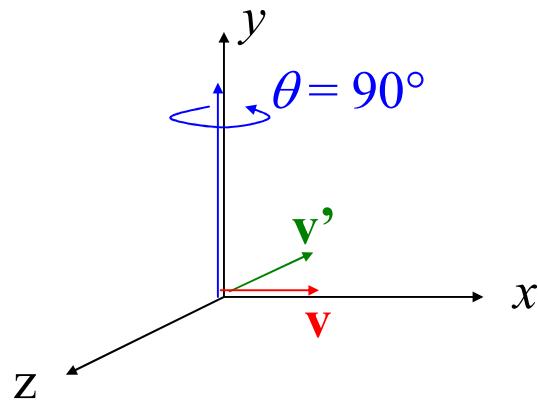
where $R = \begin{bmatrix} 1 - 2(q_2^2 + q_3^2) & 2(-q_0q_3 + q_1q_2) & 2(q_0q_2 + q_1q_3) \\ 2(q_0q_3 + q_1q_2) & 1 - 2(q_1^2 + q_3^2) & 2(-q_0q_1 + q_2q_3) \\ 2(-q_0q_2 + q_1q_3) & 2(q_0q_1 + q_2q_3) & 1 - 2(q_1^2 + q_2^2) \end{bmatrix}$

Quaternions...

- *Example*
 - A vector $\mathbf{v} = (1, 0, 0)$ is rotated about the y -axis by an angle of 90° . Determine the (unit) quaternion encoding this rotation and show that the rotated vector is $\mathbf{v}' = (0, 0, -1)$. Also find the rotation matrix corresponding to the quaternion.

Quaternions...

- *Solution*



$$\hat{\mathbf{q}} = \left[\cos\left(\frac{\theta}{2}\right), \sin\left(\frac{\theta}{2}\right)\hat{\mathbf{n}} \right] \quad \text{where } \theta=90^\circ, \hat{\mathbf{n}} = (0, 1, 0)$$

$$\Rightarrow \hat{\mathbf{q}} = \left(\frac{1}{\sqrt{2}}, 0, \frac{1}{\sqrt{2}}, 0 \right)$$

$$[0, \mathbf{v}'] = \hat{\mathbf{q}} [0, \mathbf{v}] \hat{\mathbf{q}}^{-1} = [0, \mathbf{v} \cos \theta + (\hat{\mathbf{n}} \times \mathbf{v}) \sin \theta + (\mathbf{v} \cdot \hat{\mathbf{n}}) \hat{\mathbf{n}} (1 - \cos \theta)]$$

$$= \left[0, \mathbf{0} + (0, 1, 0) \times (1, 0, 0) + \underbrace{\{(1, 0, 0) \cdot (0, 1, 0)\}}_{=0} (0, 1, 0) \right]$$

$$= [0, (0, 0, -1)]$$

$$\boxed{\therefore \mathbf{v}' = (0, 0, -1)}$$

Quaternions...

- *Solution*

$$R = \begin{bmatrix} 1 - 2(q_2^2 + q_3^2) & 2(q_1q_2 - q_0q_3) & 2(q_0q_2 + q_1q_3) \\ 2(q_0q_3 + q_1q_2) & 1 - 2(q_1^2 + q_3^2) & 2(-q_0q_1 + q_2q_3) \\ 2(-q_0q_2 + q_1q_3) & 2(q_0q_1 + q_2q_3) & 1 - 2(q_1^2 + q_2^2) \end{bmatrix}$$

$$\therefore R = \begin{bmatrix} 0 & 0 & 1 \\ 0 & 1 & 0 \\ -1 & 0 & 0 \end{bmatrix}$$

Note: $\begin{Bmatrix} v'_1 \\ v'_2 \\ v'_3 \end{Bmatrix} = \underbrace{\begin{bmatrix} 0 & 0 & 1 \\ 0 & 1 & 0 \\ -1 & 0 & 0 \end{bmatrix}}_{=R} \underbrace{\begin{Bmatrix} 1 \\ 0 \\ 0 \end{Bmatrix}}_{=\mathbf{v}} = \begin{Bmatrix} 0 \\ 0 \\ -1 \end{Bmatrix}$ (same result as before)

Quaternions...

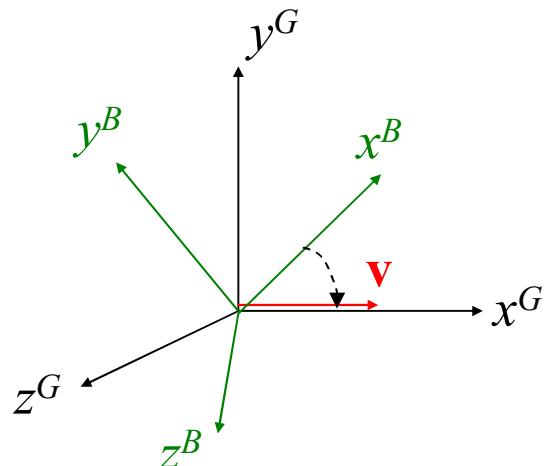
- Using quaternions to change FORs
 - Let \mathbf{q}^{GB} be the unit quaternion that rotates FOR G to FOR B.
 - Since the vector \mathbf{v} does not rotate with the FOR, it's as if \mathbf{v} rotates in the *opposite* sense of \mathbf{q}^{GB} (imagine the movement of \mathbf{v} if standing in FOR B).
 - Thus, to express \mathbf{v} in FOR B, rotate \mathbf{v}^G using the inverse quaternion $\mathbf{q}^{BG} = (\mathbf{q}^{GB})^{-1}$.

$$\begin{aligned} [0, \mathbf{v}^B] &= (\hat{\mathbf{q}}^{BG}) [0, \mathbf{v}^G] (\hat{\mathbf{q}}^{BG})^{-1} \\ &= (\hat{\mathbf{q}}^{GB})^{-1} [0, \mathbf{v}^G] (\hat{\mathbf{q}}^{GB}) \end{aligned}$$

$$\mathbf{v}^B = R^{GB} \mathbf{v}^G$$

where $R^{GB} =$

$$\begin{bmatrix} 1 - 2[(q_2^{GB})^2 + (q_3^{GB})^2] & 2(q_0^{GB} q_3^{GB} + q_1^{GB} q_2^{GB}) & 2(-q_0^{GB} q_2^{GB} + q_1^{GB} q_3^{GB}) \\ 2(q_1^{GB} q_2^{GB} - q_0^{GB} q_3^{GB}) & 1 - 2[(q_1^{GB})^2 + (q_3^{GB})^2] & 2(q_0^{GB} q_1^{GB} + q_2^{GB} q_3^{GB}) \\ 2(q_0^{GB} q_2^{GB} + q_1^{GB} q_3^{GB}) & 2(-q_0^{GB} q_1^{GB} + q_2^{GB} q_3^{GB}) & 1 - 2[(q_1^{GB})^2 + (q_2^{GB})^2] \end{bmatrix}$$

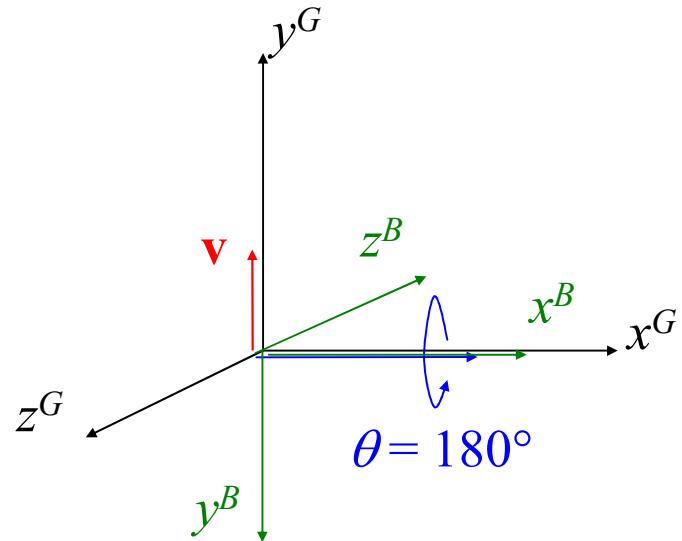


Quaternions...

- *Example*
 - A vector $\mathbf{v}^G = (0, 1, 0)$ is expressed in FOR G. FOR B is found by rotating FOR G about the vector $(1, 0, 0)^G$ (expressed in FOR G) by an angle of 180° . Determine the (unit) quaternion encoding the rotation from FOR G to FOR B. Express the vector \mathbf{v} in FOR B.

Quaternions...

- *Solution*



$$\begin{aligned}
 \hat{\mathbf{q}}^{GB} &= \left[\cos\left(\frac{\theta}{2}\right), \sin\left(\frac{\theta}{2}\right) \hat{\mathbf{n}}^G \right] = (0, 1, 0, 0) \\
 \left[0, \mathbf{v}^B \right] &= (\hat{\mathbf{q}}^{GB})^{-1} \left\{ \left[0, \mathbf{v}^G \right] (\hat{\mathbf{q}}^{GB}) \right\} \\
 &= (0, -1, 0, 0) \{ (0, 0, 1, 0) (0, 1, 0, 0) \} \\
 &= (0, -1, 0, 0) (0, 0, 0, -1) \\
 &= (0, 0, -1, 0)
 \end{aligned}$$

$$\therefore \mathbf{v}^B = (0, -1, 0)$$

Quaternions...

- Quaternion rate of change
 - consider a rotation about the body-fixed axes over a short time Δt
 - $\omega^B(t)$ is the rate of rotation about the body-fixed axes at time t
 - $\hat{\mathbf{q}}^{GB}(t)$ is the unit quaternion to go from FOR B to FOR B at time t
 - the new quaternion may be found by multiplying the old quaternion by the quaternion describing the incremental rotation

$$\hat{\mathbf{q}}^{GB}(t + \Delta t) = \left[\cos\left(\frac{|\boldsymbol{\omega}^G(t)|\Delta t}{2}\right), \frac{\boldsymbol{\omega}^G(t)}{|\boldsymbol{\omega}^G(t)|} \sin\left(\frac{|\boldsymbol{\omega}^G(t)|\Delta t}{2}\right) \right] \hat{\mathbf{q}}^{GB}(t)$$

$$\therefore \dot{\mathbf{q}}^{GB}(t) = \frac{1}{2} \hat{\mathbf{q}}^{GB}(t) [0, \boldsymbol{\omega}^B(t)]$$

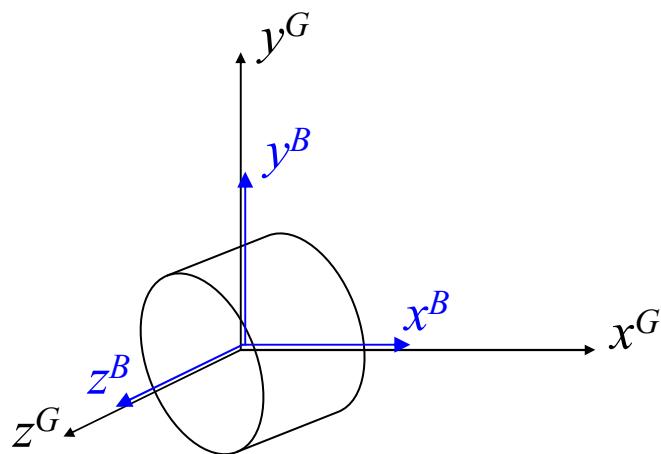
Note: $[0, \boldsymbol{\omega}^G] = (\hat{\mathbf{q}}^{GB}) [0, \boldsymbol{\omega}^B] (\hat{\mathbf{q}}^{GB})^{-1}$

- there is no singularity in the quaternion rate of change (recall that there was using Euler angles)
- only unit quaternions encode rotations so the quaternion should be normalized after it has been updated to prevent numerical “drift” resulting from numerical precision errors

Quaternions...

- *Example*

- A cylinder's body-fixed FOR (see below for the orientation of the body-fixed FOR) is found in the global FOR using the quaternion $\mathbf{q}^{GB} = (0, 1/\sqrt{2}, 0, 1/\sqrt{2})$. The cylinder rotates about its body-fixed axes with speed $\omega^B = (0, 0, 1)$. Sketch the cylinder's orientation in the global FOR and determine the rate of change of the cylinder's quaternion at this instant in time.

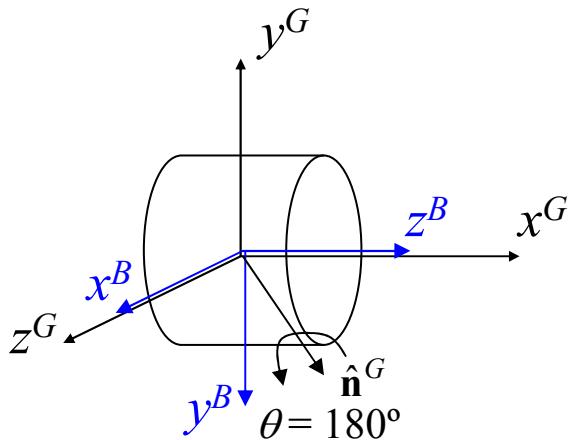


Quaternions...

- Solution*

Note: $\theta = 180^\circ$ and $\hat{\mathbf{n}}^G = (1, 0, 1)$

$$\Rightarrow \hat{\mathbf{q}} = \left(0, \frac{1}{\sqrt{2}}, 0, \frac{1}{\sqrt{2}} \right)$$



Determine the orientation of the body-fixed z-axis

$$\begin{aligned} [0, \mathbf{z}'] &= (\hat{\mathbf{q}}^{GB}) [0, (0, 0, 1)] (\hat{\mathbf{q}}^{GB})^{-1} \\ &= \left(0, \frac{1}{\sqrt{2}}, 0, \frac{1}{\sqrt{2}} \right) \left\{ (0, 0, 0, 1) \left(0, -\frac{1}{\sqrt{2}}, 0, -\frac{1}{\sqrt{2}} \right) \right\} \\ &= \left(0, \frac{1}{\sqrt{2}}, 0, \frac{1}{\sqrt{2}} \right) \left(\frac{1}{\sqrt{2}}, 0, -\frac{1}{\sqrt{2}}, 0 \right) \\ &= \left(0, \frac{1}{2} + \frac{1}{2}, 0, -\frac{1}{2} + \frac{1}{2} \right) \end{aligned}$$

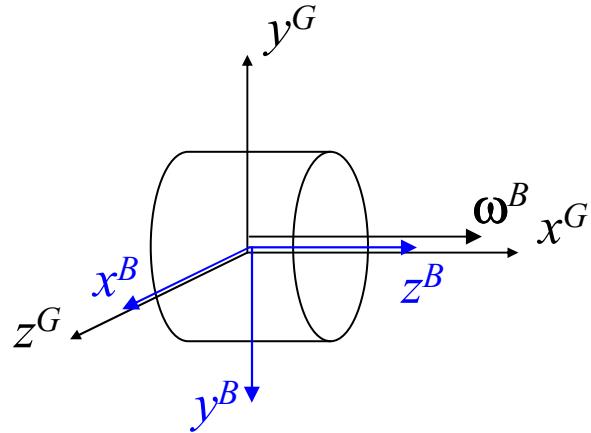
$\therefore \mathbf{z}' = (1, 0, 0)$ the body-fixed z axis points along the global x axis

Perform a similar analysis for the body-fixed x-axis

$$\begin{aligned} [0, \mathbf{x}'] &= (\hat{\mathbf{q}}^{GB}) [0, (1, 0, 0)] (\hat{\mathbf{q}}^{GB})^{-1} \\ &= \left(0, \frac{1}{\sqrt{2}}, 0, \frac{1}{\sqrt{2}} \right) \left\{ (0, 1, 0, 0) \left(0, -\frac{1}{\sqrt{2}}, 0, -\frac{1}{\sqrt{2}} \right) \right\} \\ &= \left(0, \frac{1}{\sqrt{2}}, 0, \frac{1}{\sqrt{2}} \right) \left(\frac{1}{\sqrt{2}}, 0, \frac{1}{\sqrt{2}}, 0 \right) \\ &= \left(0, \frac{1}{2} - \frac{1}{2}, 0, \frac{1}{2} + \frac{1}{2} \right) \\ \therefore \mathbf{x}' &= (0, 0, 1) \end{aligned}$$

Quaternions...

- *Solution...*



$$\begin{aligned}\dot{\mathbf{q}}^{GB} &= \left[0, \frac{1}{2} \boldsymbol{\omega}^B \right] \hat{\mathbf{q}}^{GB} \\ &= \left(0, 0, 0, \frac{1}{2} \right) \left(0, \frac{1}{\sqrt{2}}, 0, \frac{1}{\sqrt{2}} \right) \\ \therefore \dot{\mathbf{q}}^{GB} &= \left(-\frac{1}{2\sqrt{2}}, 0, \frac{1}{2\sqrt{2}}, 0 \right)\end{aligned}$$

Summary

- 3D orientations and rotations are more complex than their 2D counterparts
- Euler angles suffer from “gimbal lock” at particular angles – shouldn’t use for modeling rotational motion
- Quaternions have many advantages and are commonly used in DEM applications
 - don’t suffer from gimbal lock
 - easy to correct for numerical drift
 - multiple rotations can be represented by quaternion multiplication